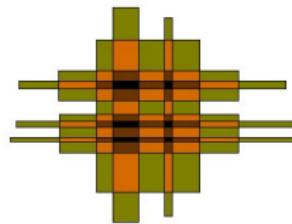


Singular integrals in Fluid Mechanics



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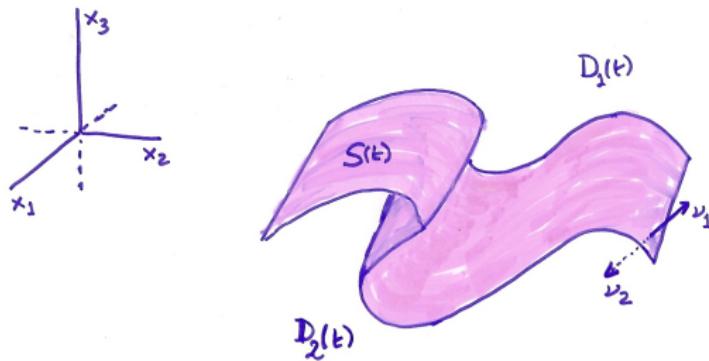
(A.C., D. CÓRDOBA AND F. GANCEDO)

- Interface evolution: the Hele-Shaw and Muskat problems. *Annals of Mathematics* **173** (2011), 1–66.
- The Raleigh-Taylor condition for the evolution of irrotational fluid interfaces. *Proc. Natl. Acad. Sci. USA* **106** (2009).
- Interface evolution: water waves in 2D. *Advances in Mathematics* **233** (2010), 120–173.
- Porous media: The Muskat problem in 3D. In press.

Bibliography

- AMBROSE, D.: Well-posedness of two-phases Darcy flow in 3-D. *Quart. Appl. Math.* **65** (2007), no 1.
- BEAR, J.: *Dynamics of Fluids in Porous Media*. Am. Elsevier, N.Y., 1972.
- CONSTANTIN, P.; PUGH, M.: Global solutions for small data to the Hele-Shaw problem. *Nonlinearity* **6** (1993).
- DARCY, H.: *Les Fontaines Publiques de la Ville de Dijon*. Ed. Dalmont, Paris, 1856.
- MUSKAT, M.: *The flow of homogeneous fluids through porous media*. N.Y., 1937.
- SIEGEL, M.; CAFLISCH, R. AND HOWISON, S.: Global existence, singular solutions and ill-posedness for the Muskat problems. *Comm. Pure and Appl. Math.* **57** (2004).

Setting



Two fluids

velocity u^j
pressure p^j
viscosity $\mu^1 \neq \mu^2$
density $\rho^1 \neq \rho^2$

- Porous media: porosity $\kappa > 0$ (homogeneous, isotropic).

- Incompressibility (mass conservation):

$$\operatorname{div}(u) = 0, \quad u = u^1 \chi_{D_1} + u^2 \chi_{D_2}$$

- Darcy's Law (conservation of momentum):

$$\frac{\mu^j}{\kappa} u^j = -\nabla p^j - (0, 0, \rho^j g) = -\nabla(p^j + g \rho^j x_3) \quad (g = \text{gravity}).$$

Immediate consequences

- $\operatorname{div}(u) = 0 \implies (u^1 - u^2) \cdot \nu_j = 0.$
- $\operatorname{curl}(u) = \Omega dS,$ Ω is a tangent vector field at $S.$
- $p^1|_S = p^2|_S.$
- \exists velocity potentials: $u^j = \nabla \phi^j$

$$\boxed{\frac{\mu^j}{\kappa} \phi^j = -p^j - \rho^j g x_3}$$

Interface motion

$$u(x) = \int_S \frac{x-z}{|x-z|^3} \wedge \Omega(z) dS(z), \quad x \notin S \quad (\text{Biôt-Savart}).$$

Taking limits in each domain:

$$u^j(x) = p.v. \int_S \frac{x-z}{|x-z|^3} \wedge \Omega(z) dS(z) + \text{tangential terms}, \quad x \in S.$$

Choosing a parametrization (isothermal coordinates)

$$X(\alpha) = X(\alpha_1, \alpha_2) \quad \begin{cases} X_{\alpha_1} \cdot X_{\alpha_2} = 0 \\ \|X_{\alpha_1}\| = \|X_{\alpha_2}\| \end{cases}$$

Two scenarios:

- Periodic in the horizontal variables;
- asymptotically flat.

Equation of motion

$$\begin{aligned}\frac{dX(\alpha)}{dt} &= p \cdot v \cdot \int \frac{X(\alpha) - X(\beta)}{|X(\alpha) - X(\beta)|^3} \wedge \omega(\alpha) d\alpha \\ &\quad + C_1(X(\alpha, t), t) X_{\alpha_1} + C_2(X(\alpha, t), t) X_{\alpha_2} \\ &= BR(X, \omega) + C_1 X_{\alpha_1} + C_2 X_{\alpha_2}.\end{aligned}$$

where

$$\text{curl}(u) = \omega(\alpha) \delta_{X-X(\alpha, t)}$$

Closing the equations: relationship between $\omega(\alpha, t)$ and $X(\alpha, t)$

Darcy's law:

$$\frac{\mu^j}{\kappa} \phi^j = -p^j - \rho^j g x_3, \quad u^j = \nabla \phi^j$$

We define

$$\Pi(\alpha) = \phi^1(X(\alpha, t), t) - \phi^2(X(\alpha, t), t).$$

Then we have

- $\omega(\alpha) = \frac{\partial \Pi}{\partial \alpha_2} X_{\alpha_1} - \frac{\partial \Pi}{\partial \alpha_1} X_{\alpha_2}$
- $\Pi(\alpha) + \frac{\mu^2 - \mu^1}{\mu^2 + \mu^1} \mathcal{D}(\Pi)(\alpha) = g \frac{\rho^2 - \rho^1}{\mu^2 + \mu^1} X_3(\alpha)$

where

$$\mathcal{D}(\Pi)(\alpha) = \int_S \frac{X(\alpha) - X(\beta)}{|X(\alpha) - X(\beta)|^3} N(\beta) \Pi(\beta) d\beta$$

and

$$N(\beta) = X_{\alpha_1} \wedge X_{\alpha_2}.$$

It is then crucial to estimate

$$\|(I + \lambda \mathcal{D})^{-1}\|, \quad |\lambda| = \left| \frac{\mu^2 - \mu^1}{\mu^2 + \mu^1} \right| < 1.$$

Equations

The problem

$$\begin{cases} \frac{dX}{dt} = BR(X, \omega) + C_1 X_{\alpha_1} + C_2 X_{\alpha_2} \\ \omega = \frac{\partial \Pi}{\partial \alpha_2} X_{\alpha_1} - \frac{\partial \Pi}{\partial \alpha_1} X_{\alpha_2} \\ \Pi(\alpha) = \left(I + \frac{\mu^2 - \mu^1}{\mu^2 + \mu^1} \mathcal{D} \right)^{-1} \left(g \frac{\rho^2 - \rho^1}{\mu^2 + \mu^1} X_3(\alpha, t) \right) \end{cases}$$

is not well-posed.

Rayleigh-Taylor:

$$\begin{aligned} \sigma(\alpha, t) &= (\nabla p^2 - \nabla p^1) \cdot (\nu^2 - \nu^1) \\ &= (\mu^2 - \mu^1) BR(X, \omega) \cdot (X_{\alpha_1} \wedge X_{\alpha_2}) + (\rho^2 - \rho^1) (X_{\alpha_1} \wedge X_{\alpha_2})_3 \\ &> 0 \end{aligned}$$



Geometric constraint: not self-intersecting surface

$$F(x) = \sup_{\alpha \neq \beta} \frac{\|\alpha - \beta\|}{\|X(\alpha) - X(\beta)\|} < \infty.$$

Theorem

If

- $X_0(\alpha) \in H^4(\mathbb{R}^2)$;
- $\sigma(\alpha, 0) > 0$;
- $F(X_0) < \infty$;

then $\exists T > 0$ and a unique solution $X \in C^1([0, T]; H^4)$.

Energy estimates

$$\begin{aligned}\frac{d}{dt} \|X\|_{H^k}^2 &\leq P(\|X\|_{H^k} + \|F(X)\|_{L^\infty}) \\ &\quad - \sum_{j=0}^k \frac{2^{3/2}}{\mu^1 + \mu^2} \int \frac{\sigma(\alpha)}{|\nabla X(\alpha)|^3} \frac{\partial^{k-j}}{\partial \alpha_1} \frac{\partial^j}{\partial \alpha_2} X(\alpha) \\ &\quad \cdot \Lambda \left(\frac{\partial^{k-j}}{\partial \alpha_1} \frac{\partial^j}{\partial \alpha_2} X(\alpha) \right) d\alpha\end{aligned}$$

where

$$\boxed{\Lambda = \sqrt{-\Delta}}$$

Then we use the pointwise inequality

$$\theta(x) \cdot \Lambda\theta(x) \geq \frac{1}{2}\Lambda\theta^2(x)$$

to get a polynomial estimate in the norm H^k .

However, we need to control the evolution of all the involved quantities, namely,

- $\sigma > 0$.
- $F < \infty$.
- C_j (isothermal).
- $\|I + \lambda\mathcal{D}\|^{-1}, \quad |\lambda| < 1$.